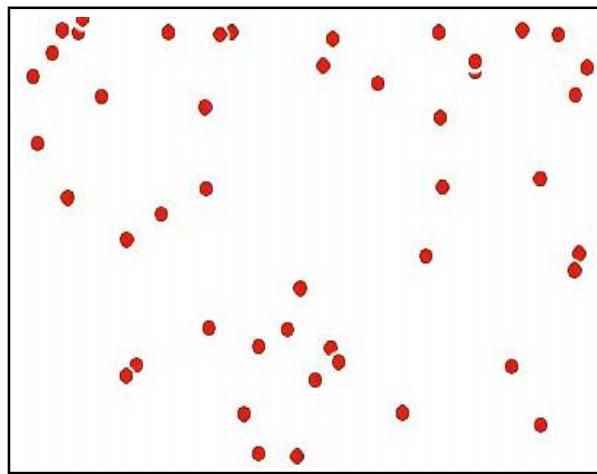


Du microscopique au macroscopique

Description microscopique (Position, vitesse)



Propriétés macroscopiques (Variables d'état)

Pression (Pa) Température (K)

$$PV=nRT$$

Volume (m^3)

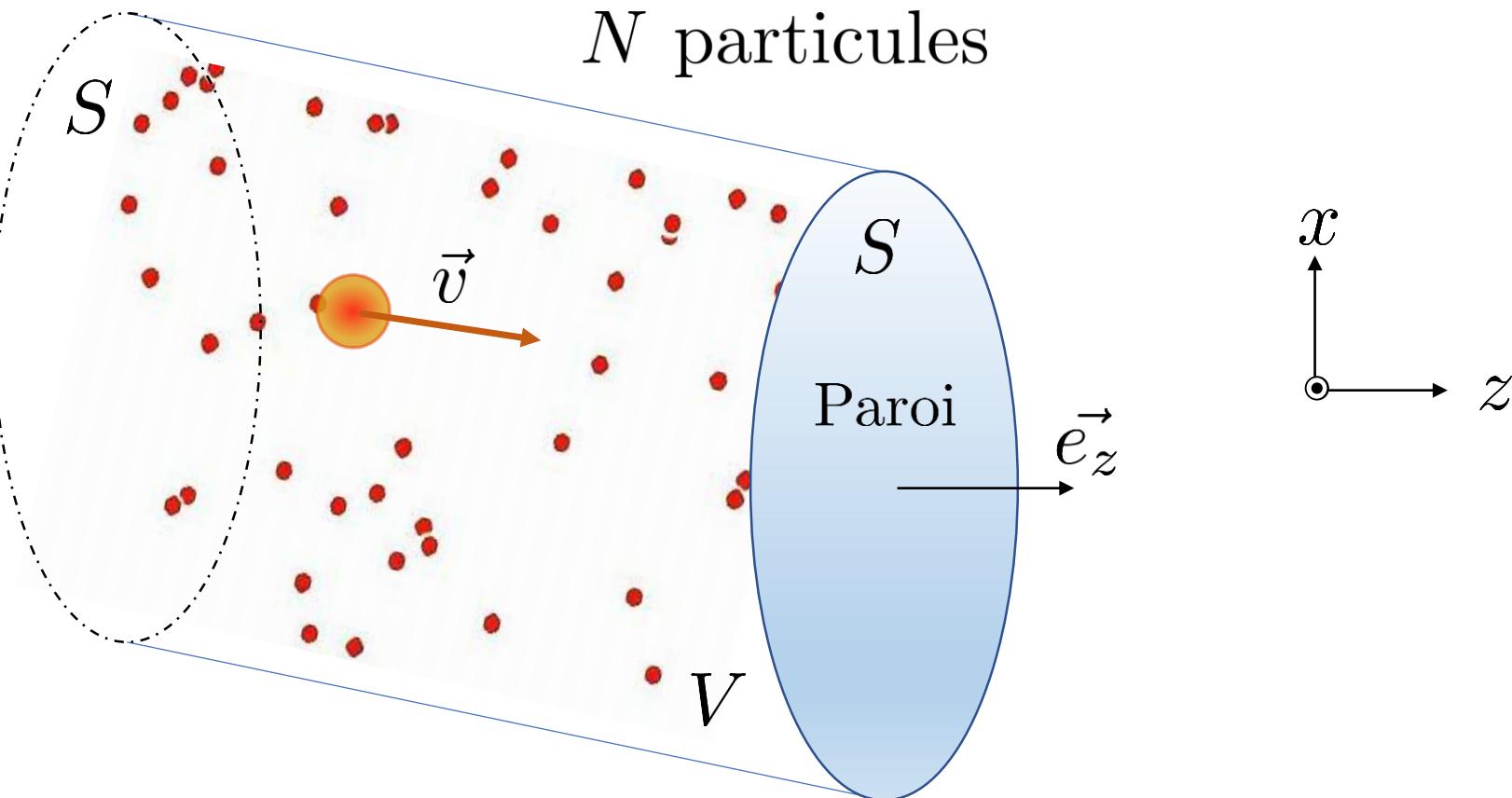
Qte de matière (mol)

Loi des gaz parfaits

Problème compliqué :

1mm³ de gaz $\sim 10^{16}$ particules

Système thermodynamique considéré

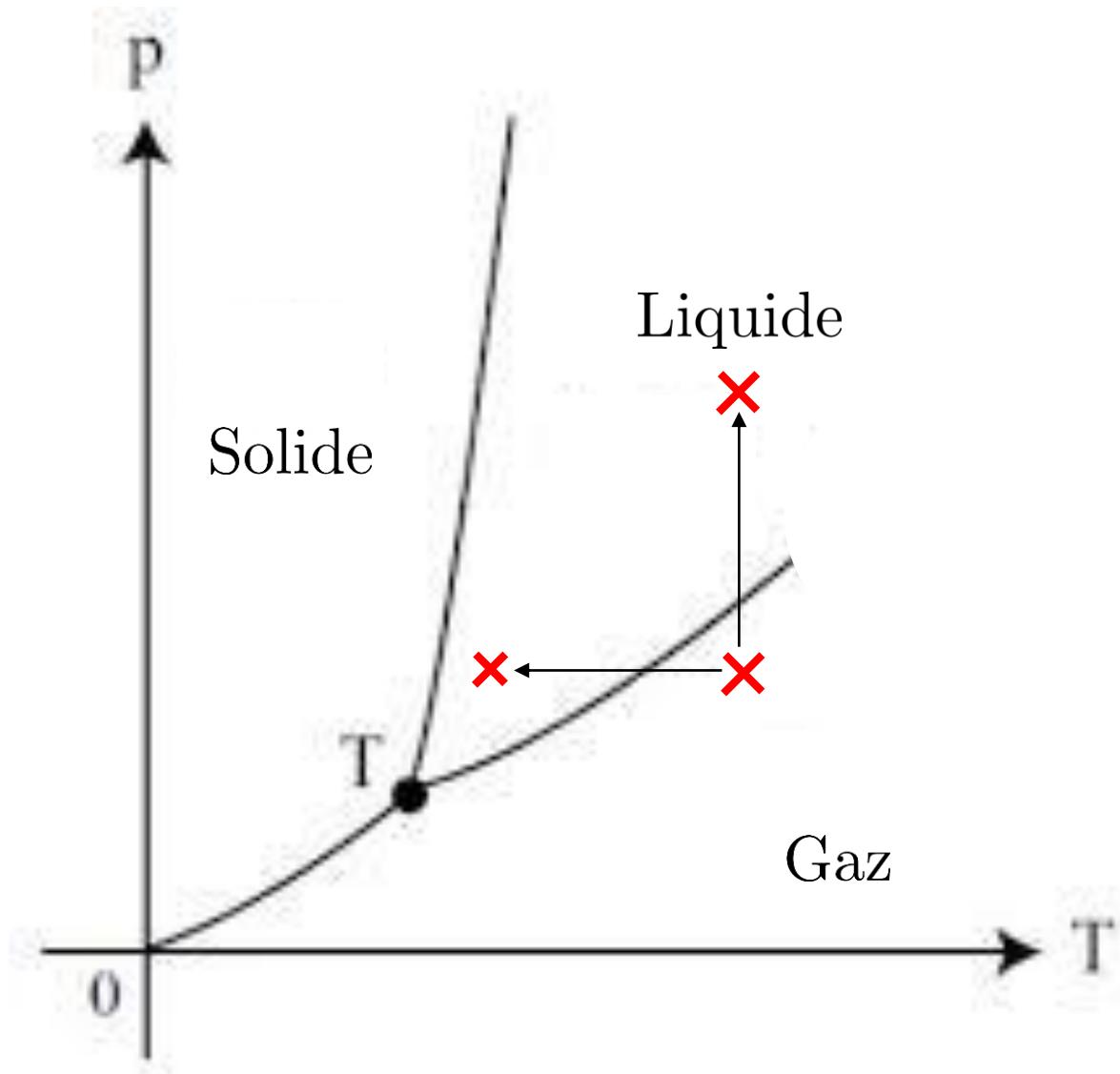


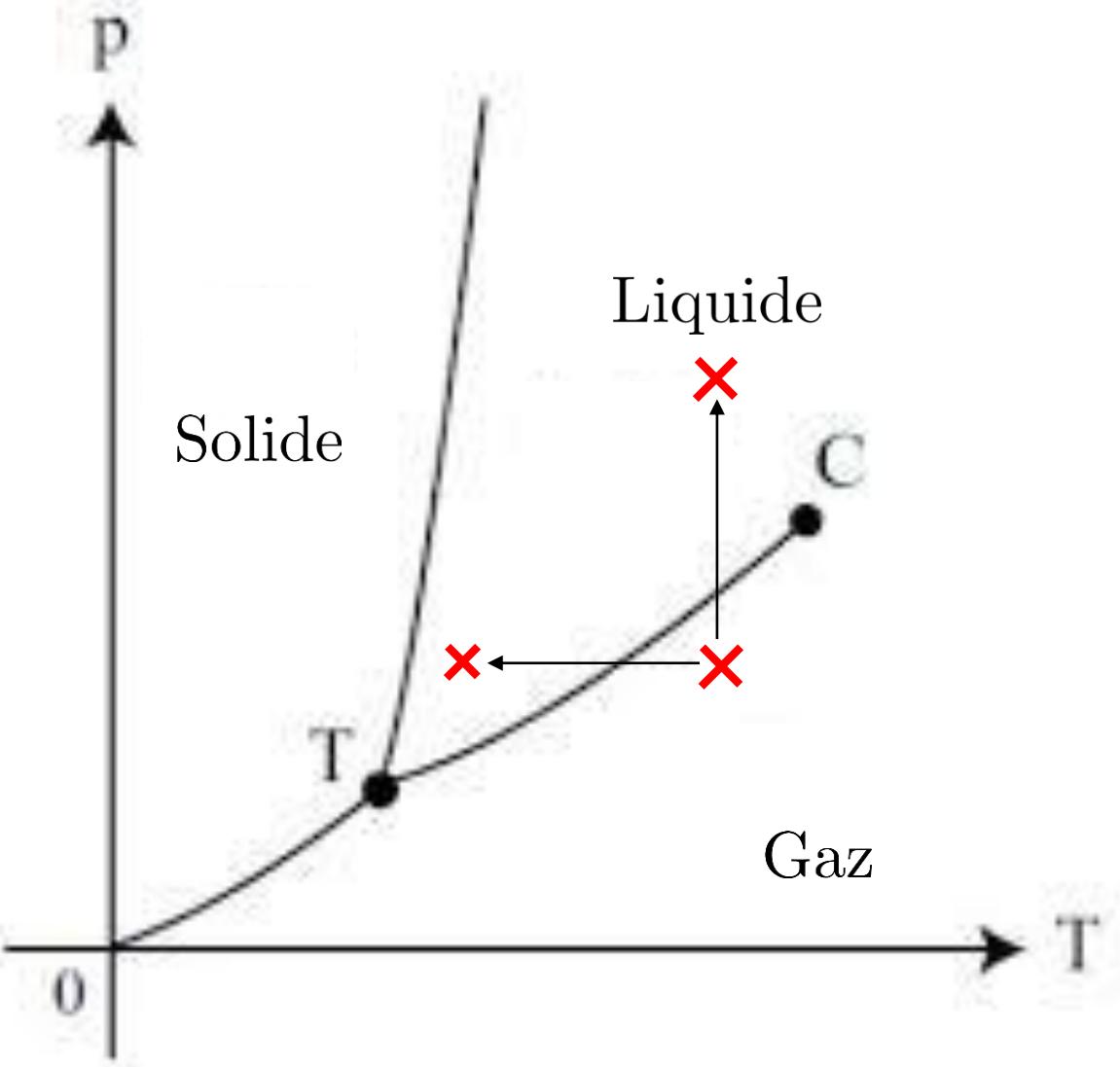
Distribution des vitesses (Maxwell Boltzmann)

$$dP(v) = \left(\frac{m}{2\pi k_B T}\right)^{3/2} e^{-\frac{mv^2}{k_B T}} dv$$

Résultat admis

Probabilité que $v \in [v, v + dv]$





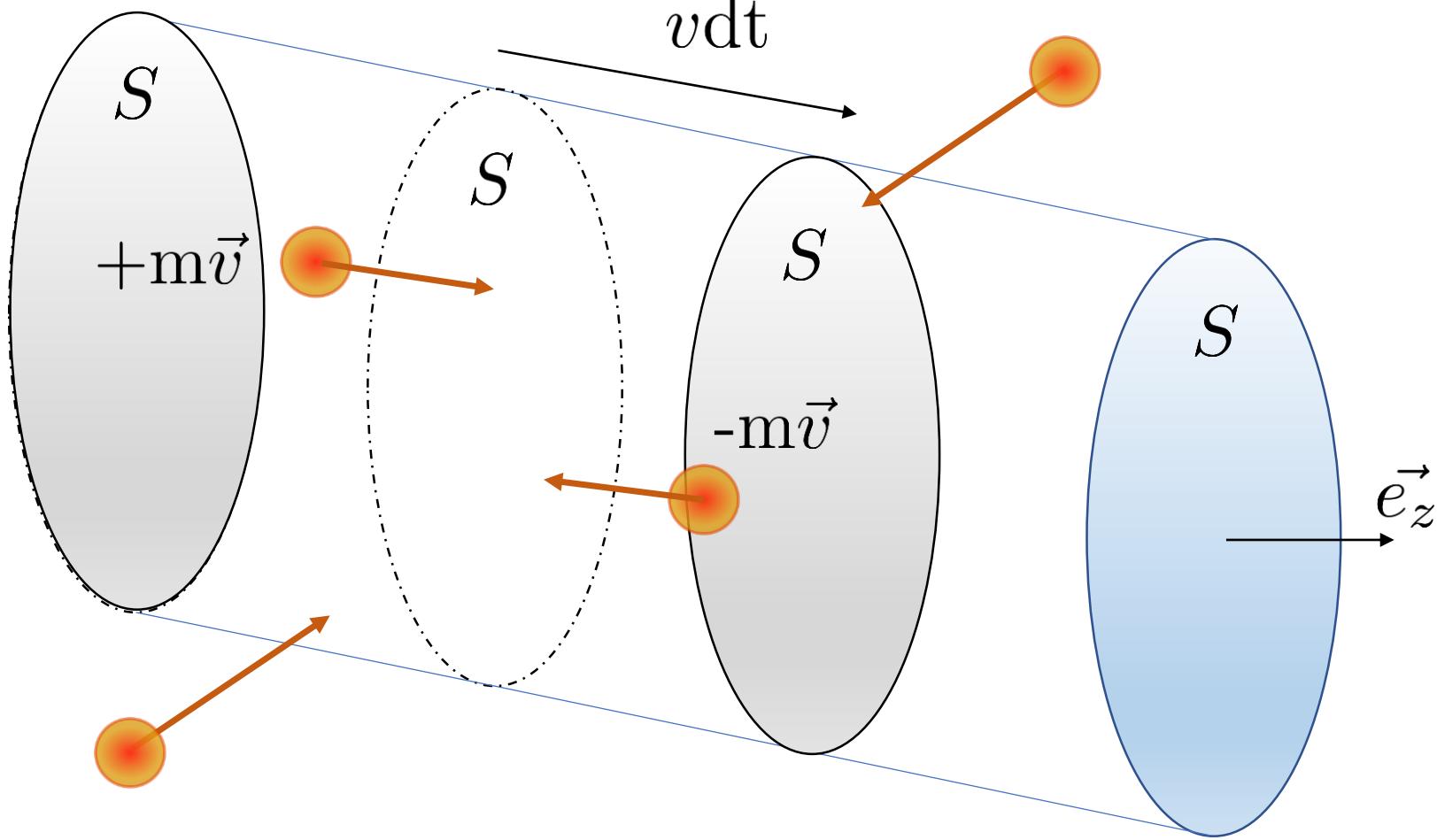


Merci de votre
attention

Monsieur, pourquoi les atomes qui rentrent
par les surfaces latérales ne contribuent pas ?

x

z



Beattie-Bridgeman Equation of State

$$P = \frac{R_u T}{\bar{v}^2} \left(1 - \frac{c}{\bar{v} T^3} \right) (\bar{v} + B) - \frac{A}{\bar{v}^2}$$

The constants appearing in the above equation are given in Table 3–4 for various substances. The Beattie-Bridgeman equation is known to be reasonably accurate for densities up to about $0.8\rho_{\text{cr}}$, where ρ_{cr} is the density of the substance at the critical point.

Benedict-Webb-Rubin Equation of State

$$P = \frac{R_u T}{\bar{v}} + \left(B_0 R_u T - A_0 - \frac{C_0}{T^2} \right) \frac{1}{\bar{v}^2} + \frac{b R_u T - a}{\bar{v}^3} + \frac{a\alpha}{\bar{v}^6} + \frac{c}{\bar{v}^3 T^2} \left(1 + \frac{\gamma}{\bar{v}^2} \right) e^{-\gamma/\bar{v}^2}$$

The values of the constants appearing in this equation are given in Table 3–4. This equation can handle substances at densities up to about $2.5\rho_{\text{cr}}$. In 1962, Strobridge further extended this equation by raising the number of constants to 16 (Fig. 3–59).

Distribution de Maxwell Boltzmann des vitesses

$$dP(v) = \left(\frac{m}{2\pi k_B T}\right)^{3/2} e^{-\frac{mv^2}{k_B T}} dv$$

$$\langle v_z^2 \rangle = \int_{-\infty}^{+\infty} v_z^2 P(v) dv = \frac{k_B T}{m}$$

$$\langle v^2 \rangle = \int_{-\infty}^{+\infty} v^2 P(v) dv = \frac{3k_B T}{m}$$